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I. Numerical Differentiation near the Limits of a Difference Table.

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1. *Introduction.*

THIS paper has its immediate origin in the implications of a remark by Bradfield and Southwell⁽¹⁾. From the values of a function calculated (approximately) for eleven equally spaced arguments they wished to calculate derivatives of the function for the same arguments, using the known finite difference formulae. They remark that "the order of the highest derivative which can be calculated is different at different sections"—implying that this order is that of the highest difference which can be used. If the finite difference formulae are restricted to those containing only forward, backward, or central differences respectively, the implication is correct—but there must exist formulae of a "mixed"

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type which enable the whole ten differences to be utilized at any argument, *i. e.*, formulæ commencing as central difference formulæ, switching over to forward (or backward) differences when the edge of the difference table is reached*. The route through the difference table would then be $\longrightarrow \nearrow$ or $\longrightarrow \searrow$. Similar considerations apply near the beginning or end of any difference table. We are not aware that the necessary coefficients have hitherto been published, and the object of this paper is to supply them.

In section 2 the main results, in the form of tables of coefficients, are described, together with (it is hoped sufficient) discussion and examples to enable those who so wish to use the tables without mastering or reading the underlying theory. In the last section the theory, calculation, and checking of the formulæ are considered.

Numerical differentiation—by *any* formula—is of necessity relatively inaccurate when the values of the function are not exact (*e. g.*, rounded-off or experimental), and the smaller the tabular interval the greater is the relative effect of such inexactitude. It is therefore desirable to use a large tabular interval, even though this entails the use of differences of highish order; we give coefficients for differences up to the twelfth. This inaccuracy increases as one proceeds to higher derivatives, when it is usually combined with a progressive loss of significant figures; we have therefore refrained from giving complete tables for derivatives of higher order than the fourth, although they have been computed for all orders up to the twelfth. Should they be needed, they can be constructed by applying the rules developed in section 3.3 to the entries of the basic Tables I. and IV.

The possible and/or probable errors in the derivatives calculated from the "mixed" type formulæ vary with the argument. A discussion of them is desirable, and is in hand.

2. The Tables and the Method of Use.

2.1. *Notation.*—The usual † notations for forward, backward, and central differences, as exhibited in the following schemes, will be used;

* Formulæ starting diagonally, and finishing horizontally, will also exist. They are less convenient and more laborious in use, and will not be considered here.

† These differ from astronomical usage (Comrie⁽²⁾); there forward, backward, and central differences are not distinguished by different symbols, Δ being used for central differences, while mean differences are identifiable by difference of parity between index and doubled suffix.

the actual numbers in the same relative positions will, of course, be identical.

x	Forward.	Backward.	Central.
a	y_0	y_0	y_0
$a+w$	Δy_0	∇y_1	$\delta y_{\frac{1}{2}}$
$a+2w$	$\Delta^2 y_0$	$\nabla^2 y_2$	$\delta^2 y_1$
$a+3w$	$\Delta^3 y_0$	$\nabla^3 y_3$	$\delta^3 y_{\frac{3}{2}}$
$a+4w$	$\Delta^4 y_0$	$\nabla^4 y_4$	$\delta^4 y_2$

along with $\mu\delta^n y_m = \frac{1}{2}(\delta^n y_{m+\frac{1}{2}} + \delta^n y_{m-\frac{1}{2}})$.

The tabular interval is denoted by w . The differential operator will be denoted by D , and we shall use the abbreviation

$$(D^q y)_{x=a+mw} = D^q y_m, \text{ or sometimes } D_m^q.$$

We abbreviate $\delta^n y_m$ similarly to δ_m^n .

2.2. Description of the Tables.—Table I. (p. 11) gives the Gregory-Newton coefficients $N_{q,n}$ in the formulæ

$$w^q D^q y_m = \begin{cases} \sum_{n=q}^{\infty} (-1)^{n-q} N_{q,n} \Delta^n y_m \\ \sum_{n=q}^{\infty} N_{q,n} \nabla^n y_m \end{cases}$$

for $q=0(1)12$ and $n \geq 12$.

Table II. (p. 12) gives the corresponding central difference coefficients $C_{q,n}$ in the formulæ

$$w^q D^q y_m = \begin{cases} \sum_{r=0}^{\infty} (-1)^r C_{q,q+2r} \delta^{q+2r} y_m & (q \text{ even}) \\ \sum_{r=0}^{\infty} (-1)^r C_{q,q+2r} \mu \delta^{q+2r} y_m & (q \text{ odd}) \end{cases}$$

Tables III., 1-4 (facing p. 12), give the coefficients in the mixed type formulæ for the first four derivatives at the tabulated values of the argument. Their arrangement is such that the coefficient to be used as a multiplier occupies the same relative position as the difference which it is to multiply. The "route" is indicated by the arrows; central difference coefficients are given in bolder type; the central difference coefficient (always unity) on the left of the table for D^q multiplies $\delta^q y_0$ if q is even and $\mu\delta^q y_0$ if q is odd. In order further to clarify the mode of usage of the tables, illustrative examples are worked out in detail (see pp. 4-6).

Tables IV. (p. 13) and V. (p. 14), for arguments at mid-interval, correspond to Tables I. and II., and give $N'_{q,n}$ and $C'_{q,n}$ for the formulæ

$$w^q D^q y_{m-\frac{1}{2}} = \begin{cases} \sum_{n=q}^{\infty} (-)^{n-q} N'_{q,n} \Delta^n y_m \\ \sum_{n=q}^{\infty} N'_{q,n} \nabla^n y_{m-1} \\ \sum_{r=0}^{\infty} (-)^r C'_{q,q+2r} \mu \delta^{q+2r} y_{m-\frac{1}{2}} \quad (q \text{ even}) \\ \sum_{r=0}^{\infty} (-)^r C'_{q,q+2r} \delta^{q+2r} y_{m-\frac{1}{2}} \quad (q \text{ odd}) \end{cases}$$

Tables VI., 0-3 (facing p. 14), give the coefficients in the mixed type formulæ for the function and its first three derivatives at mid-interval.

2.3. *Numerical Examples.*—As material we employ values of the function $\cos x$ (x in radians) at interval 0.1, to ten decimals. The values and differences are given on p. 5.

We will use these differences, and the coefficients from Tables III. 1, III. 2, and VI. 1, to calculate $y'(0.2)$, $y''(0.2)$, and $y'(0.25)$ respectively.

(i) $y'(0.2)$

Coefficient (Table III., 1)	×	difference (p. 5)	=	term in result
+1		$\mu \delta_2 =$		-1983 38381
-1/6		$\mu \delta_2^2 = 19$		3 30288 7
+1/30		$\Delta_0^5 = \delta_{2\frac{1}{2}}^5 = -$		24686
-1/60		$\Delta_0^6 = \delta_3^6 = -$		9533
+1/105		$\Delta_0^7 = \delta_{3\frac{1}{2}}^7 = +$		347
-1/168		$\Delta_0^8 = \delta_4^8 = +$		81
				<hr/>
				-1986 69330 9

Since $w=0.1$, this means that

$$y'(0.2) = -0.19866 9331.$$

The exact value is $-\sin 0.2 = -0.19866 93307 95 \dots$

(ii) $y''(0.2)$

Coefficient (Table III., 2)	×	difference (p. 5)	=	term in result
+1		$\delta_2^2 =$		-979 25012
-1/12		$\delta_2^4 = +9$		78432
+1/90		$\Delta_0^6 = \delta_3^6 = -$		9533
-1/90		$\Delta_0^7 = \delta_{3\frac{1}{2}}^7 = +$		347
+47/5040		$\Delta_0^8 = \delta_4^8 = +$		81
-19/2520		$\Delta_0^9 = \delta_{4\frac{1}{2}}^9 = +$		25
				<hr/>
				-980 06657 2

x	$\cos x$	δ	δ^2	δ^3	δ^4	δ^5	δ^6	δ^7	δ^8	δ^9
0.0	1.00000	00000								
0.1	0.99500	41653	-994 17528							
0.2	0.98006	65778	-979 25012	+14 92510	+9 78432					
0.3	0.95533	64891	-954 54064	+24 70948	+9 53746	-24686				
0.4	0.92106	09940	-920 29370	+34 24604	+9 19527	-34219	-9533	+347	+ 81	+25
0.5	0.87758	25619	-876 85149	+43 44221	+8 76122	-43405	-9186	+428	+106	-59
0.6	0.82533	56149	-824 64806	+52 20343	+8 23959	-52103	-8758	+534	+ 47	
0.7	0.76484	21873	-764 20504	+60 41302	+7 63572	-60387	-8224	+581		
0.8	0.69670	67093	-696 12630	+68 07874	+6 95542	-68030	-7643			
0.9	0.62160	99683	-621 09214	+75 03416						
1.0	0.54030	23059	-540 30							

The terms in bold type are mean central differences.

Since $w^2=0.01$, this means that

$$y''(0.2) = -0.98006\ 657.$$

This may be compared with the true value,

$$-\cos 0.2 = -0.98006\ 65778\dots$$

(iii) $y'(0.25)$

Coefficient (Table VI., 1)	×	difference (p. 5)	=	term in result
+1		$\delta_{2\frac{1}{2}} =$		-2473 00887
-1/24		$\delta_{2\frac{1}{2}}^3 = +24\ 70948$		- 1 02956 2
+3/640		$\delta_{2\frac{1}{2}}^5 = - 24686$		- 115 7
-5/7168		$\Delta_0^7 = \delta_{3\frac{1}{2}}^7 = + 347$		- 2
+5/7168		$\Delta_0^8 = \delta_{\frac{1}{2}}^8 = + 81$		+ 1
				-2474 03959 0

Thus $y'(0.25) = -0.24740\ 3959$,

while $-\sin 0.25 = -0.24740\ 39592\ 5\dots$

Scrutiny of (i) and (ii) will show how and why decimals, and hence significant figures, are progressively lost as we reduce the interval or as we calculate higher derivatives, whenever the tabular interval is less than unity.

3. Derivation of Formulæ and Calculation of Coefficients.

3.1. *Gregory-Newton Coefficients.*—Denoting as usual by E the operation of adding w to the argument, or its equivalent of increasing any suffix by unity,

$$y_1 = Ey_0 = (1 + \Delta)y_0 = e^{wD}y_0$$

where the last expression is the symbolic form of Taylor's series. Thus

$$E = (1 + \Delta) = e^{wD}$$

so that

$$wD = \log(1 + \Delta)$$

$$= \Delta - \frac{1}{2}\Delta^2 + \frac{1}{3}\Delta^3 - \frac{1}{4}\Delta^4 + \dots$$

Consequently

$$\begin{aligned} w^q D^q y_m &= (\Delta - \frac{1}{2}\Delta^2 + \frac{1}{3}\Delta^3 - \frac{1}{4}\Delta^4 + \dots)^q y_m \\ &= \sum_{n=q}^{\infty} (-1)^{n-q} N_{q,n} \Delta^n y_m \end{aligned}$$

Similarly

$$E^{-1} = (1 - \nabla) = e^{-wD}$$

and

$$wD = -\log(1 - \nabla)$$

$$= \nabla + \frac{1}{2}\nabla^2 + \frac{1}{3}\nabla^3 + \frac{1}{4}\nabla^4 + \dots$$

so that
$$w^q D^q y_m = (\nabla + \frac{1}{2} \nabla^2 + \frac{1}{3} \nabla^3 + \frac{1}{4} \nabla^4 + \dots)^q y_m$$

$$= \sum_{n=q}^{\infty} N_{q,n} \nabla^n y_m$$

The coefficients $N_{q,n}$ can thus be calculated by involution of the logarithmic series. Other—and less laborious—methods exist. They depend on the Gregory-Newton interpolation formula

$$y(a+rw) = y_r = \sum \frac{r(r-1)(r-2) \dots (r-n+1)}{n!} \Delta^n y_0$$

Differentiating q times with respect to r gives

$$w^q D^q y_r = \sum \frac{d^q \{r(r-1)(r-2) \dots (r-n+1)\} / dr^q}{n!} \Delta^n y_0$$

so that

$$n! N_{q,n} = (-)^{n-q} [d^q \{r(r-1)(r-2) \dots (r-n+1)\} / dr^q]_{r=0}$$

$$= [d^q \{r(r+1)(r+2) \dots (r+n-1)\} / dr^q]_{r=0}$$

which is essentially Markoff's formula. The values of the derivatives when $r=0$ can evidently be found as simple multiples of the coefficients, and this leads to a second method of calculating $N_{q,n}$.

From the last result we can also derive a third method. Thus

$$(n+1)! N_{q,n+1} = [(d/dr)^q \{r(r+1)(r+2) \dots (r+n)\}]_{r=0}$$

$$= [(r+n)(d/dr)^q \{r(r+1) \dots (r+n-1)\}]_{r=0}$$

$$+ q(d/dr)^{q-1} \{r(r+1)(r+2) \dots (r+n-1)\}_{r=0}$$

or, on reduction,

$$(n+1)N_{q,n+1} = nN_{q,n} + qN_{q-1,n}$$

This recurrence formula, coupled with the fact that $N_{q,q} = 1$, enables Table I. to be built up column by column, and provides what is probably the quickest and least laborious way of doing so. Actually all three methods have been employed, and they check one another.

It has seemed worth while also to provide formulæ for mid-interval, since in numerical processes the time often comes when a reduction of the interval is imperative, and the simplest method of so doing is to halve it. The corresponding coefficients are obtained from formulæ very similar to those above.

Putting $r = s - \frac{1}{2}$, the Gregory-Newton formula becomes

$$y_{s-\frac{1}{2}} = \sum_{n=0}^{\infty} \frac{(s-\frac{1}{2})(s-\frac{3}{2}) \dots (s-n+\frac{1}{2})}{n!} \Delta^n y_0$$

whence we derive, as in the case of $N_{q,n}$

$$n! N'_{q,n} = [(d/ds)^q \{(s+\frac{1}{2})(s+\frac{3}{2}) \dots (s+n-\frac{1}{2})\}]_{s=0}$$

and

$$2(n+1)N'_{q,n+1} = (2n+1)N'_{q,n} + 2qN'_{q-1,n}$$

This recurrence formula, coupled with $N'_{q,q}=1$, $N'_{q,0}=0$ ($q>0$), enables Table IV. to be constructed. From the above it is also evident that the N' are multiples of the coefficients in the polynomial

$$(x+1)(x+3)\dots(x+2n-1),$$

which provides a second method of calculating them. A third method depends on the fact that

$$w^q D^q y_{-\frac{1}{2}} = (\Delta - \frac{1}{2}\Delta^2 + \frac{1}{3}\Delta^3 - \dots)^q (1 + \Delta)^{-\frac{1}{2}} y_0$$

3.2. *Central Difference Formulæ.*—These have been given, up to the terms involving the twelfth difference, by Comrie (2). It will transpire in our next subsection that the coefficients automatically emerge in the process of calculating the coefficients for the "mixed" formulæ. To calculate them independently we have

$$\delta = E^{\frac{1}{2}} - E^{-\frac{1}{2}} = 2 \sinh \frac{1}{2} wD$$

so that

$$wD = 2 \sinh^{-1} (\delta/2)$$

Also

$$\mu = \frac{1}{2}(E^{\frac{1}{2}} + E^{-\frac{1}{2}}) = \cosh \frac{1}{2} wD = \sqrt{1 + \frac{1}{4}\delta^2}$$

Consequently

$$w^q D^q = \begin{cases} \{2 \sinh^{-1}(\delta/2)\}^q & \begin{cases} q \text{ even, interval points} \\ q \text{ odd, mid-interval} \end{cases} \\ \mu \{2 \sinh^{-1}(\delta/2)\}^q (1 + \frac{1}{4}\delta^2)^{-\frac{1}{2}} & \begin{cases} q \text{ odd, interval points} \\ q \text{ even, mid-interval} \end{cases} \end{cases}$$

For further information and details concerning differentiation formulæ see Steffensen (4) §7, and §18.

3.3. *Mixed Formulæ.*—Although the general rule by which the coefficients in the "mixed" formulæ are derived is the same in all cases, its deduction will probably be more easily followed if we analyse the particular case of the first derivative.

We have

$$\begin{aligned} wDy_1 &= (\Delta - \frac{1}{2}\Delta^2 + \frac{1}{3}\Delta^3 - \frac{1}{4}\Delta^4 \dots)y_1 \\ &= \Delta y_1 - (\frac{1}{2}\Delta^2 - \frac{1}{3}\Delta^3 + \frac{1}{4}\Delta^4 \dots)(1 + \Delta)y_0 \\ &= \Delta y_1 - \frac{1}{2}\Delta^2 y_0 + \{(\frac{1}{2} - \frac{1}{3})\Delta^3 - (\frac{1}{4} - \frac{1}{3})\Delta^4 \dots\}y_0 \\ &= \mu\delta y_1 - \frac{1}{6}\Delta^3 y_0 + \frac{1}{12}\Delta^4 y_0 - \dots \end{aligned}$$

since

$$\mu\delta y_1 = \frac{1}{2}(\Delta y_1 + \Delta y_0) = \Delta y_1 - \frac{1}{2}\Delta^2 y_0$$

Similarly

$$\begin{aligned} wDy_2 &= \mu\delta y_2 - (\frac{1}{6}\Delta^3 - \frac{1}{12}\Delta^4 + \frac{1}{20}\Delta^5 \dots)y_1 \\ &= \mu\delta y_2 - \frac{1}{6}\Delta^3 y_1 + (\frac{1}{12}\Delta^4 - \frac{1}{20}\Delta^5 \dots)(1 + \Delta)y_0 \\ &= \mu\delta y_2 - \frac{1}{6}\mu\delta^3 y_1 + \{(\frac{1}{12} - \frac{1}{20})\Delta^5 - (\frac{1}{20} - \frac{1}{30})\Delta^6 + \dots\}y_0 \end{aligned}$$

The process can be indefinitely continued, for the law of formation of the coefficients is now clear. If the accompanying table represents the coefficients, with

1	d_2			
	c_1	d_3		
1	c_2	d_4		
	b_1	c_3	d_4	
1	b_2	c_4		
	a_1	b_3	c_4	
	a_2	b_4		
		a_3	b_4	
			a_4	

$$a_n = (-)^n N_{a, n+q}$$

then $b_n = a_{n-1} + a_n$

$$c_n = b_{n-1} + b_n$$

and so on. The upward progress in any column ceases when the corresponding backward Gregory-Newton

coefficient is reached. The reproduction of the known central difference coefficients, and later of the backward Gregory-Newton coefficients, provides a check (practically conclusive and complete) upon the accuracy of the calculations.

Modifications are still necessary to convert a table so constructed into the corresponding member of the set III. or VI. In the case of odd derivatives at tabular points, or of even derivatives at mid-interval—that is, where the central difference formula involves mean differences—the central line of coefficients in the original table is repeated three times, at interval $\frac{1}{2}w$, with $\delta_{-\frac{1}{2}}$, $\mu\delta_0$ and $\delta_{\frac{1}{2}}$ respectively in III., 1, 3, or with δ_0 , $\mu\delta_{\frac{1}{2}}$ and δ_1 in VI., 0, 2. In the case of even derivatives at tabular points, or of odd derivatives at mid-interval, the original table has a central line of zeros, and the entries in the lines above and below this are equal. In Tables III. 2, 4, these entries are repeated three times, with δ_{-1} , δ_0 , and δ_1 respectively (or with $\delta_{-\frac{1}{2}}$, $\delta_{\frac{1}{2}}$ and $\delta_{1\frac{1}{2}}$ in Tables VI. 1, 3), instead of twice, at interval w ; the zeros are not printed in the tables.

The suffixes also need consideration. In Tables III. the sum of the outermost pair in any column is zero. In Tables VI., however, an asymmetry appears, owing to the fact that the tables relate to suffix $+\frac{1}{2}$; the sum of the outermost suffixes in any column is, in this case, unity.

It may also be mentioned that the number of 1's in the left-hand column *as computed* is equal to the order of the derivative, for the interval point schemes, but one more than this for mid-interval schemes.

To illustrate these points we reproduce a few columns of the tables for D^2y and D^3y as computed; these should be compared with the corresponding portions of Tables III. 2, and III. 3.

The tables could also have been exhibited in the form of "hexagon" diagrams (Fraser⁽³⁾), and portions corresponding to parts of the tables just mentioned are given, to exhibit this possibility. In the hexagon diagrams horizontal lines indicate addition or subtraction, sloping lines indicate multiplication, and any continuous route from left to right may be followed.

D^2y_0

Original Table.						Hexagon Diagram.				
				137/180						
		5/6								
	1	11/12	-1/12	-13/180						
1		-1/12	0	1/90						
1	0	-1/12	0	1/90						
	-1	11/12	1/12	-13/180						
			-5/6							
				137/180						

D^3y_0

				29/15						
		15/8								
		7/4		7/120						
	1	3/2	1/8	-1/15						
		1/2	-1/8	7/120						
1		-1/4	1/8	-1/15						
1		-3/2	-1/8	7/120						
		7/4								
			-15/8							
				29/15						

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The authors wish to record, firstly their thanks to the Royal Society for a grant in aid of the publication of this paper, and secondly their appreciation of the care and courtesy with which the printers have endeavoured to meet their wishes, especially as regards the printing and arrangement of the tables.

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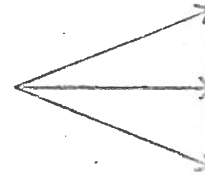
TABLE I. $N_{q,n}$

q	n	1	2	3	4	5	6	7	8	9	10	11	12	q
0	1	1	0	0	0	0	0	0	0	0	0	0	0	0
1	1	1	1	1	1	1	1	1	1	1	1	1	1	1
1	2	1	2	3	4	5	6	7	8	9	10	11	12	1
1	3	1	3	6	10	15	21	28	36	45	55	67	81	1
1	4	1	4	10	20	35	56	84	120	165	220	287	368	1
1	5	1	5	15	35	70	126	210	330	495	715	1001	1365	1
1	6	1	6	21	56	126	252	462	756	1155	1716	2431	3327	1
1	7	1	7	28	84	210	462	924	1638	2772	4557	7141	10583	1
1	8	1	8	36	126	315	756	1638	3360	5775	9724	15181	22780	1
1	9	1	9	45	165	420	1050	2430	5040	9009	14784	23185	34968	1
1	10	1	10	55	210	546	1365	3360	7560	13650	23184	38166	57656	1
1	11	1	11	67	264	735	1911	4620	10920	21420	38166	63523	103776	1
1	12	1	12	81	336	945	2520	6435	16380	33642	61416	105831	181548	1
1	13	1	13	97	420	1287	3432	8910	22680	47520	91416	163548	284296	1
1	14	1	14	115	528	1638	4410	11550	30030	63523	124638	231841	414288	1
1	15	1	15	137	660	2163	5775	14940	39690	84450	163548	300301	537966	1
1	16	1	16	161	828	2925	7770	20448	52920	112860	231841	431286	789168	1
1	17	1	17	197	1035	3864	10296	27462	71416	151815	316236	598548	1111296	1
1	18	1	18	245	1344	5005	13608	36430	97245	207360	431286	829168	1565280	1
1	19	1	19	307	1764	6720	18009	47520	124638	270270	564480	1163548	2228496	1
1	20	1	20	385	2340	9180	24312	63523	163548	349680	731286	1518168	3003024	1
1	21	1	21	481	3120	12540	33006	86430	226800	475200	1001280	2073600	4312896	1
1	22	1	22	597	4158	17160	44736	115500	300300	635232	1360320	2842968	5985480	1
1	23	1	23	745	5511	23640	61416	163548	414288	884160	1815480	3816624	8003040	1
1	24	1	24	937	7254	32460	84450	226800	500304	1058304	2228496	4641600	9724800	1
1	25	1	25	1177	9528	44460	115500	300300	672000	1414284	2925000	6141600	12844800	1
1	26	1	26	1477	12510	60480	157416	414288	900300	1915200	4003040	8441600	17652000	1
1	27	1	27	1841	16380	82800	216036	564480	1246380	2644800	5441600	11412800	23841600	1
1	28	1	28	2385	21630	112860	292500	756000	1635480	3496800	7312800	15181680	31623600	1
1	29	1	29	3033	28800	151815	396900	1029600	2268000	4752000	10012800	20736000	43128960	1
1	30	1	30	3913	38640	204480	544500	1414284	3162360	6720000	14142840	29250000	61416000	1
1	31	1	31	5077	52800	274620	714164	1815480	4142880	8841600	18154800	38166240	80030400	1
1	32	1	32	6677	72540	364300	972450	2431200	5003040	10583040	22284960	46416000	97248000	1
1	33	1	33	8861	99450	500304	1360320	3496800	7312800	15181680	31623600	66416000	141428400	1
1	34	1	34	11801	134400	714164	1815480	4641600	9724800	20736000	43128960	91416000	191520000	1
1	35	1	35	15777	180090	1029600	2644800	6141600	12844800	27027000	56448000	116354800	243120000	1
1	36	1	36	21161	243120	1494000	3643000	8441600	18154800	39690000	84416000	176520000	370236000	1
1	37	1	37	28457	324600	2044800	5003040	11550000	24312000	50030400	105830400	222849600	464160000	1
1	38	1	38	38413	431280	2841600	7003040	16354800	34968000	73128000	151816800	316236000	664160000	1
1	39	1	39	51857	577500	3864000	9724500	22680000	47520000	100128000	207360000	431289600	914160000	1
1	40	1	40	70481	777000	5292000	13603200	30030000	63523000	136032000	284296800	598548000	1246380000	1
1	41	1	41	95981	1058400	7312800	18154800	40030400	84416000	176520000	370236000	789168000	1635480000	1
1	42	1	42	131401	1414284	9945000	24312000	50030400	105830400	222849600	464160000	972480000	2073600000	1
1	43	1	43	180777	1915200	13603200	31623600	66416000	141428400	292500000	614160000	1284480000	2702700000	1
1	44	1	44	247777	2541600	18154800	41428800	88416000	181548000	381662400	800304000	1635480000	3496800000	1
1	45	1	45	339137	3496800	24312000	50030400	105830400	222849600	464160000	972480000	2073600000	4312896000	1
1	46	1	46	464881	4752000	33006000	70030400	141428400	292500000	614160000	1284480000	2702700000	5644800000	1
1	47	1	47	635237	6352300	44460000	97245000	204480000	431289600	900304000	1815480000	3816624000	8003040000	1
1	48	1	48	864481	8441600	60480000	136032000	284296800	614160000	1284480000	2702700000	5644800000	11635480000	1
1	49	1	49	1169777	11412840	82800000	181548000	396900000	844160000	1765200000	3702360000	7891680000	16354800000	1
1	50	1	50	1584137	15181680	112860000	243120000	500304000	1058304000	2228496000	4641600000	9724800000	20736000000	1
1	51	1	51	2147537	20736000	151816800	330060000	690304000	1414284000	2925000000	6141600000	12844800000	27027000000	1
1	52	1	52	2919937	28416000	204480000	444600000	944160000	1915480000	4003040000	8441600000	17652000000	37023600000	1
1	53	1	53	3961337	38640000	284160000	614288000	1284480000	2702700000	5644800000	11635480000	24312000000	50030400000	1
1	54	1	54	5351737	51428400	386400000	844160000	1815480000	3816624000	8003040000	16354800000	34968000000	73128000000	1
1	55	1	55	7262137	69030400	529200000	1155000000	2431200000	5003040000	10583040000	22284960000	46416000000	97248000000	1
1	56	1	56	9872537	94416000	731280000	1635480000	3496800000	7312800000	15181680000	31623600000	66416000000	141428400000	1
1	57	1	57	13382937	128448000	994500000	2160320000	4641600000	9724800000	20736000000	43128960000	91416000000	191520000000	1
1	58	1	58	18113337	171600000	1360320000	2925000000	6141600000	12844800000	27027000000	56448000000	116354800000	243120000000	1
1	59	1	59	24463737	231280000	1815480000	4003040000	8441600000	18154800000	38166240000	80030400000	163548000000	349680000000	1
1	60	1	60	33114137	316236000	2431200000	5003040000	10583040000	22284960000	46416000000	97248000000	207360000000	431289600000	1
1	61	1	61	44764537	424160000	3300600000	7003040000	14142840000	29250000000	61416000000	128448000000	270270000000	564480000000	1
1	62	1	62	60414937	577500000	4446000000	9724500000	20448000000	43128960000	90030400000	181548000000	381662400000	800304000000	1
1	63	1	63	82065337	777000000	6048000000	13603200000	28429680000	61416000000	128448000000	270270000000	564480000000	1163548000000	1
1	64	1	64	110715737	1058400000	8280000000	18154800000	39690000000	84416000000	176520000000	370236000000	789168000000	1635480000000	1
1	65	1	65	148466137	1414284000	11286000000	24312000000	50030400000	105830400000	222849600000	464160000000	972480000000	2073600000000	1
1	66	1	66	198216537	1815480000	13603200000	29250000000	61416000000	128448000000	270270000000	564480000000	1163548000000	2431200000000	1
1	67	1	67	263966937	2431200000	18154800000	40030400000	84416000000	181548000000	381662400000	800304000000	1635480000000	3496800000000	1
1	68	1	68	351717337	3162360000	24312000000	50030400000	105830400000	222849600000	464160000000	972480000000	2073600000000	4312896000000	1
1	69	1	69	468467737	4241600000	33006000000	7							

TABLE II. $C_{q,n}$

q	n	0	1	2	3	4	5	6	7	8	9	10	11	12	q
0		1													0
1			1		$\frac{1}{6}$	0	$\frac{1}{30}$	$\frac{1}{90}$	$\frac{1}{140}$	0	$\frac{1}{630}$	0	$\frac{1}{2772}$	0	1
2				1		$\frac{1}{12}$		$\frac{1}{90}$	$\frac{1}{560}$	$\frac{1}{560}$	$\frac{1}{3024}$	$\frac{1}{3150}$	$\frac{1}{16632}$		2
3					1		$\frac{1}{4}$	$\frac{7}{120}$	$\frac{7}{120}$	$\frac{7}{240}$	$\frac{41}{3024}$	$\frac{41}{7560}$	$\frac{479}{151200}$		3
4						1		$\frac{1}{6}$		$\frac{7}{240}$		$\frac{41}{7560}$	$\frac{479}{453600}$		4
5							1		$\frac{1}{3}$		$\frac{13}{144}$	$\frac{13}{240}$	$\frac{139}{6048}$		5
6								1		$\frac{1}{4}$		$\frac{13}{240}$	$\frac{139}{12096}$		6
7									1		$\frac{5}{12}$		$\frac{31}{240}$		7
8										1		$\frac{1}{3}$	$\frac{31}{360}$		8
9											1		$\frac{1}{2}$		9
10												1	$\frac{5}{12}$		10
11														1	11
12															12

$$10^q D^q y_m = \begin{cases} \sum_{r=0}^{\infty} (-1)^r C_{q,q+2r} \delta^{q+2r} y_m & (q \text{ even}) \\ \sum_{r=0}^{\infty} (-1)^r C_{q,q+2r} \mu \delta^{q+2r} y_m & (q \text{ odd}) \end{cases}$$



[To face p. 12.]

TABLE III. 1. Dy_0

$$\begin{aligned}
& \delta_{-1}^2 + \frac{1}{3} \delta_{-1}^3 + \frac{1}{6} \delta_{-1}^3 + \frac{1}{6} \mu \delta_0^3 + \frac{1}{3} \delta_{1\frac{1}{2}}^3 + \frac{1}{12} \delta_1^4 + \frac{1}{3} \delta_{1\frac{1}{2}}^4 - \frac{1}{4} \delta_2^4 + \frac{1}{5} \delta_{2\frac{1}{2}}^5 + \frac{1}{20} \delta_{1\frac{1}{2}}^5 + \frac{1}{30} \mu \delta_0^5 + \frac{1}{30} \delta_{\frac{1}{2}}^5 + \frac{1}{12} \delta_1^6 + \frac{1}{20} \delta_{1\frac{1}{2}}^6 - \frac{1}{4} \delta_2^6 + \frac{1}{5} \delta_{2\frac{1}{2}}^6 - \frac{1}{6} \delta_3^6 + \frac{1}{7} \delta_{3\frac{1}{2}}^7 + \frac{1}{105} \delta_{1\frac{1}{2}}^7 + \frac{1}{140} \mu \delta_0^7 + \frac{1}{140} \delta_{\frac{1}{2}}^7 + \frac{1}{105} \delta_{1\frac{1}{2}}^8 + \frac{1}{168} \delta_2^8 - \frac{1}{168} \delta_{2\frac{1}{2}}^8 + \frac{1}{56} \delta_3^8 + \frac{1}{7} \delta_{3\frac{1}{2}}^8 - \frac{1}{8} \delta_4^8 + \frac{1}{9} \delta_{4\frac{1}{2}}^9 + \frac{1}{90} \delta_{-5}^{10} + \frac{1}{90} \delta_{-4}^{10} + \frac{1}{360} \delta_{-3}^{10} - \frac{1}{840} \delta_{-2}^{10} + \frac{1}{1260} \delta_{-1}^{10} - \frac{1}{2772} \mu \delta_0^{10} - \frac{1}{2772} \delta_{\frac{1}{2}}^{10} - \frac{1}{1260} \delta_1^{10} + \frac{1}{840} \delta_2^{10} - \frac{1}{360} \delta_3^{10} + \frac{1}{90} \delta_4^{10} - \frac{1}{10} \delta_5^{10} + \frac{1}{11} \delta_{5\frac{1}{2}}^{11} - \frac{1}{110} \delta_{-4\frac{1}{2}}^{11} + \frac{1}{495} \delta_{-3\frac{1}{2}}^{11} - \frac{1}{1320} \delta_{-2\frac{1}{2}}^{11} + \frac{1}{2310} \delta_{-1\frac{1}{2}}^{11} - \frac{1}{2772} \mu \delta_0^{11} - \frac{1}{2772} \delta_{\frac{1}{2}}^{11} + \frac{1}{2310} \delta_{1\frac{1}{2}}^{11} - \frac{1}{1320} \delta_{2\frac{1}{2}}^{11} + \frac{1}{495} \delta_{3\frac{1}{2}}^{11} - \frac{1}{110} \delta_{4\frac{1}{2}}^{11} + \frac{1}{11} \delta_{5\frac{1}{2}}^{11} - \frac{1}{132} \delta_{-5}^{12} - \frac{1}{132} \delta_{-5}^{12} + \frac{1}{660} \delta_{-4}^{12} - \frac{1}{1980} \delta_{-3}^{12} + \frac{1}{3960} \delta_{-2}^{12} - \frac{1}{5544} \delta_{-1}^{12} + \frac{1}{5544} \delta_{1}^{12} - \frac{1}{3960} \delta_{2}^{12} + \frac{1}{1980} \delta_{3}^{12} - \frac{1}{660} \delta_{4}^{12} + \frac{1}{132} \delta_{5}^{12} - \frac{1}{12} \delta_0^{12}
\end{aligned}$$

TABLE III. 2. $D^2 y_0$

$$\begin{aligned}
 & + 1 \delta_{-1\frac{1}{2}} \\
 & + \frac{11}{12} \delta_{-2} \\
 & + \frac{5}{6} \delta_{-2\frac{1}{2}} \\
 & + \frac{137}{180} \delta_{-3} \\
 & + \frac{7}{10} \delta_{-3\frac{1}{2}} \\
 & + \frac{1}{12} \delta_{-4} \\
 & - \frac{1}{12} \delta_{-4} \\
 & - \frac{1}{12} \delta_{-1} \\
 & - \frac{1}{12} \delta_0 \\
 & - \frac{1}{12} \delta_1 \\
 & + \frac{11}{12} \delta_2 \\
 & - \frac{5}{6} \delta_{2\frac{1}{2}} \\
 & + \frac{137}{180} \delta_3 \\
 & + \frac{1}{90} \delta_{-1} \\
 & + \frac{1}{90} \delta_0 \\
 & + \frac{1}{90} \delta_1 \\
 & + \frac{1}{12} \delta_{1\frac{1}{2}} \\
 & - \frac{13}{180} \delta_2 \\
 & - \frac{5}{6} \delta_{2\frac{1}{2}} \\
 & + \frac{137}{180} \delta_3 \\
 & + \frac{1}{90} \delta_{-1} \\
 & + \frac{1}{90} \delta_0 \\
 & + \frac{1}{90} \delta_1 \\
 & - \frac{1}{90} \delta_{1\frac{1}{2}} \\
 & + \frac{47}{5040} \delta_2 \\
 & + \frac{11}{180} \delta_{2\frac{1}{2}} \\
 & + \frac{137}{180} \delta_3 \\
 & - \frac{7}{10} \delta_{3\frac{1}{2}} \\
 & + \frac{363}{560} \delta_{-4} \\
 & - \frac{29}{560} \delta_{-3} \\
 & + \frac{47}{5040} \delta_{-2} \\
 & - \frac{1}{560} \delta_{-1} \\
 & + \frac{1}{560} \delta_0 \\
 & - \frac{1}{560} \delta_1 \\
 & + \frac{47}{5040} \delta_2 \\
 & - \frac{29}{560} \delta_3 \\
 & + \frac{363}{500} \delta_{\frac{1}{2}} \\
 & + \frac{761}{1260} \delta_{\frac{3}{2}} \\
 & - \frac{223}{5040} \delta_{\frac{5}{2}} \\
 & + \frac{19}{2520} \delta_{\frac{7}{2}} \\
 & - \frac{1}{560} \delta_{\frac{9}{2}} \\
 & + \frac{1}{3150} \delta_{\frac{10}{2}} \\
 & + \frac{1}{3150} \delta_{\frac{10}{1}} \\
 & - \frac{37}{25200} \delta_{\frac{10}{2}} \\
 & + \frac{17}{2800} \delta_{\frac{10}{2}} \\
 & - \frac{481}{12600} \delta_{\frac{10}{\frac{1}{2}}} \\
 & + \frac{7129}{12600} \delta_{\frac{10}{5}} \\
 & + \frac{761}{1260} \delta_{-4\frac{1}{2}} \\
 & - \frac{223}{5040} \delta_{-3\frac{1}{2}} \\
 & + \frac{19}{2520} \delta_{-2\frac{1}{2}} \\
 & - \frac{1}{560} \delta_{-1\frac{1}{2}} \\
 & + \frac{1}{3150} \delta_{-1} \\
 & + \frac{1}{3150} \delta_0 \\
 & + \frac{1}{3150} \delta_1 \\
 & - \frac{37}{25200} \delta_{\frac{10}{2}} \\
 & + \frac{17}{2800} \delta_{\frac{10}{2}} \\
 & - \frac{481}{12600} \delta_{\frac{10}{\frac{1}{2}}} \\
 & + \frac{7129}{12600} \delta_{\frac{10}{5}} \\
 & + \frac{671}{1260} \delta_{-5\frac{1}{2}} \\
 & - \frac{419}{12600} \delta_{-4\frac{1}{2}} \\
 & + \frac{31}{6300} \delta_{-3\frac{1}{2}} \\
 & - \frac{29}{25200} \delta_{-2\frac{1}{2}} \\
 & + \frac{1}{3150} \delta_{-1\frac{1}{2}} \\
 & - \frac{1}{3150} \delta_{\frac{11}{1\frac{1}{2}}} \\
 & + \frac{29}{25200} \delta_{\frac{11}{2\frac{1}{2}}} \\
 & - \frac{31}{6300} \delta_{\frac{11}{3\frac{1}{2}}} \\
 & + \frac{419}{12600} \delta_{\frac{11}{4\frac{1}{2}}} \\
 & - \frac{671}{1260} \delta_{\frac{11}{5\frac{1}{2}}} \\
 & - \frac{4861}{12600} \delta_{-5} \\
 & + \frac{3349}{8 \cdot 31600} \delta_{-4} \\
 & - \frac{743}{8 \cdot 31600} \delta_{-3} \\
 & + \frac{107}{4 \cdot 15800} \delta_{-2} \\
 & - \frac{1}{16632} \delta_{-1} \\
 & - \frac{1}{16632} \delta_0 \\
 & - \frac{1}{16632} \delta_1 \\
 & + \frac{107}{4 \cdot 15800} \delta_2 \\
 & - \frac{743}{8 \cdot 31600} \delta_3 \\
 & + \frac{3349}{8 \cdot 31600} \delta_4 \\
 & - \frac{4861}{12600} \delta_5 \\
 & + \frac{83711}{1 \cdot 66320} \delta_{-6} \\
 & - \frac{4861}{1 \cdot 66320} \delta_{-5} \\
 & + \frac{3349}{8 \cdot 31600} \delta_{-4} \\
 & - \frac{743}{8 \cdot 31600} \delta_{-3} \\
 & + \frac{107}{4 \cdot 15800} \delta_{-2} \\
 & - \frac{1}{16632} \delta_{-1} \\
 & - \frac{1}{16632} \delta_0 \\
 & - \frac{1}{16632} \delta_1 \\
 & + \frac{107}{4 \cdot 15800} \delta_2 \\
 & - \frac{743}{8 \cdot 31600} \delta_3 \\
 & + \frac{3349}{8 \cdot 31600} \delta_4 \\
 & - \frac{4861}{12600} \delta_5 \\
 & + \frac{83711}{1 \cdot 66320} \delta_6
 \end{aligned}$$

TABLE IV. $N'_{q,n}$

q	n	0	1	2	3	4	5	6	7	8	9	10	11	12
0	→	1	1	2	5	128	63	231	429	6435	12155	46189	88179	6 76039
1	→	1	3	8	23	11	563	1627	88060	1423	15 93269	77 59469	317 30711	465 22243
2	→	1	1	1	3	43	95	12139	25333	81227	4 08233	1215 63409	2461 83839	3 28081 17961
3	→	1	1	1	2	24	48	5760	11320	35840	2 15040	516 09090	1032 19200	1 36249 34400
4	→	1	1	1	1	1	23	29	8197	2317	51 42611	11 11619	3161 11237	5905 69373
5	→	1	1	1	1	1	5	101	287	14861	12103	107 49419	245 63869	65980 23581
6	→	1	1	1	1	1	2	24	48	1920	1280	9 67680	19 35360	4644 86400
7	→	1	1	1	1	1	1	3	24	130	14927	19627	41 24677	8 29385
8	→	1	1	1	1	1	1	1	1	6	1152	1152	1 93536	32256
9	→	1	1	1	1	1	1	1	1	8	213	39209	1 10539	370 06861
10	→	1	1	1	1	1	1	1	1	24	16	1920	3840	9 67680
11	→	1	1	1	1	1	1	1	1	4	233	445	58933	44359
12	→	1	1	1	1	1	1	1	1	1	24	24	1920	900
														2 55731
														5700
														131
														4
														119
														24
														2
														1
														6
														1

1790

2549
2550

2551
2552

$$N'_{q,q} = 1$$

$$N'_{0,n} = 1, 3, \dots, (2n-1)/2, \dots, 2n$$

$${}_{q^0}D^q y_{m-1} = \sum_{n=q}^{\infty} (-)^{n-q} N'_{q,n} \Delta^n y_m = \sum_{n=q}^{\infty} N'_{q,n} \nabla^n y_{m-1}$$

$$2(n+1)N'_{q,n+1} = (2n+1)N'_{q,n} + 2qN'_{q-1,n}$$

14 Numerical Differentiation near the Limits of a Difference Table.

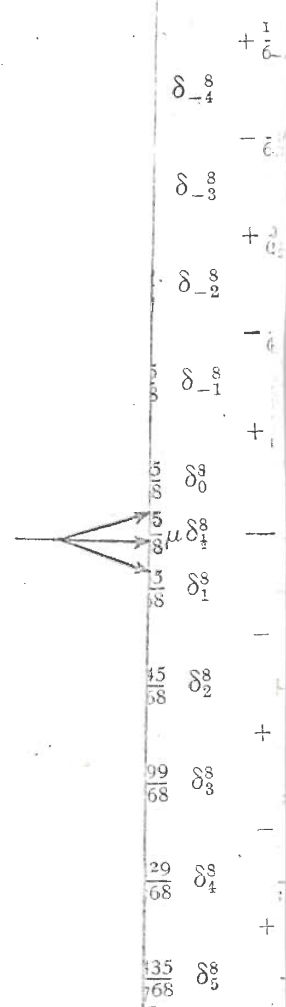
TABLE V.

q	n	0	1	2	3	4	5	6	7	8	9	10	11	12
0	1	1												
1	1		$\frac{1}{24}$											
2	1			$\frac{1}{24}$										
3	1				1									
4	1					1								
5	1						1							
6	1							1						
7	1								1					
8	1									1				
9	1										1			
10	1											1		
11	1												1	
12	1													1

q	0	1	2	3	4	5	6	7	8	9	10	11	12
0													
1													
2													
3													
4													
5													
6													
7													
8													
9													
10													
11													
12													

Handwritten notes in red ink:
 $2553 \rightarrow$
 $2554 \rightarrow$
 $2555 \rightarrow$

$$w^q D^q y_{m-\frac{1}{2}} = \begin{cases} \sum_{r=0}^{\infty} (-1)^r C_{q, q+2r}^r \delta^{q+2r} \mu^q y_{m-\frac{1}{2}} & (q \text{ even}) \\ \sum_{r=0}^{\infty} (-1)^r C_{q, q+2r}^r \delta^{q+2r} y_{m-\frac{1}{2}} & (q \text{ odd}) \end{cases}$$



BICKLEY

TABLE VI. 1. Dy_1

$$\begin{aligned}
& \frac{1027}{1020} \delta_{-3}^6 + \frac{88069}{107520} \delta_{-3\frac{1}{2}}^7 + \frac{1423}{1792} \delta_{-4}^8 + \frac{15}{20} \frac{93269}{64384} \delta_{-4\frac{1}{2}}^9 + \frac{77}{103} \frac{59469}{21920} \delta_{-5}^{10} + \frac{317}{432} \frac{30711}{53760} \delta_{-5\frac{1}{2}}^{11} + \frac{465}{648} \frac{22243}{80640} \delta_{-6}^{12} \\
& - \frac{3043}{107520} \delta_{-2\frac{1}{2}}^7 - \frac{2689}{107520} \delta_{-3}^8 - \frac{46027}{20} \frac{46027}{64384} \delta_{-3\frac{1}{2}}^9 - \frac{51719}{25} \frac{51719}{80480} \delta_{-4}^{10} - \frac{164}{9083} \frac{88341}{28960} \delta_{-4\frac{1}{2}}^{11} - \frac{21}{1297} \frac{47647}{61280} \delta_{-5}^{12} \\
& \frac{11}{960} \delta_{-2}^6 + \frac{143}{35840} \delta_{-1\frac{1}{2}}^7 + \frac{59}{17920} \delta_{-2}^8 - \frac{1195}{20} \frac{1195}{64384} \delta_{-1\frac{1}{2}}^9 - \frac{475}{10} \frac{475}{32192} \delta_{-2}^{10} + \frac{2513}{259} \frac{2513}{52256} \delta_{-1\frac{1}{2}}^{11} + \frac{973}{129} \frac{973}{76128} \delta_{-2}^{12} \\
& \frac{3}{1440} \delta_{-1}^6 - \frac{5}{7168} \delta_{-1\frac{1}{2}}^7 - \frac{5}{7168} \delta_{-1}^8 + \frac{35}{2} \frac{35}{94912} \delta_{-1\frac{1}{2}}^9 + \frac{35}{2} \frac{35}{94912} \delta_{-1}^{10} - \frac{63}{28} \frac{63}{83584} \delta_{-1\frac{1}{2}}^{11} - \frac{63}{28} \frac{63}{83584} \delta_{-1}^{12} \\
& \begin{matrix} \rightarrow \\ \rightarrow \\ \rightarrow \end{matrix} - \frac{5}{7168} \delta_{\frac{1}{2}}^7 \begin{matrix} \rightarrow \\ \rightarrow \\ \rightarrow \end{matrix} + \frac{35}{2} \frac{35}{94912} \delta_{\frac{1}{2}}^9 \begin{matrix} \rightarrow \\ \rightarrow \\ \rightarrow \end{matrix} - \frac{63}{28} \frac{63}{83584} \delta_{\frac{1}{2}}^{11} \\
& \frac{3}{960} \delta_2^6 + \frac{143}{35840} \delta_{2\frac{1}{2}}^7 + \frac{59}{17920} \delta_3^8 - \frac{1195}{20} \frac{1195}{64384} \delta_{2\frac{1}{2}}^9 - \frac{475}{10} \frac{475}{32192} \delta_3^{10} + \frac{2513}{259} \frac{2513}{52256} \delta_{2\frac{1}{2}}^{11} + \frac{973}{129} \frac{973}{76128} \delta_3^{12} \\
& \frac{1027}{1020} \delta_1^6 + \frac{88069}{107520} \delta_{1\frac{1}{2}}^7 + \frac{1423}{1792} \delta_2^8 + \frac{15}{20} \frac{93269}{64384} \delta_{1\frac{1}{2}}^9 + \frac{77}{103} \frac{59469}{21920} \delta_2^{10} + \frac{317}{432} \frac{30711}{53760} \delta_{1\frac{1}{2}}^{11} + \frac{465}{648} \frac{22243}{80640} \delta_1^{12} \\
& - \frac{3043}{107520} \delta_{\frac{3}{2}}^7 - \frac{2689}{107520} \delta_{\frac{4}{2}}^8 - \frac{46027}{20} \frac{46027}{64384} \delta_{\frac{3}{2}}^9 - \frac{51719}{25} \frac{51719}{80480} \delta_{\frac{5}{2}}^{10} - \frac{164}{9083} \frac{88341}{28960} \delta_{\frac{3}{2}}^{11} - \frac{21}{1297} \frac{47647}{61280} \delta_{\frac{6}{2}}^{12} \\
& + \frac{88069}{107520} \delta_{\frac{7}{2}}^7 - \frac{1423}{1792} \delta_{\frac{5}{2}}^8 + \frac{15}{20} \frac{93269}{64384} \delta_{\frac{9}{2}}^9 - \frac{77}{103} \frac{59469}{21920} \delta_{\frac{10}{2}}^{10} + \frac{317}{432} \frac{30711}{53760} \delta_{\frac{11}{2}}^{11} - \frac{465}{648} \frac{22243}{80640} \delta_{\frac{12}{2}}^{12}
\end{aligned}$$

TABLE VI. 3. D^3y_4

$$\begin{aligned}
 & + \frac{5005}{774} \frac{69373}{14400} \delta_{-6}^{12} \\
 & + \frac{3161}{516} \frac{11237}{09600} \delta_{-5\frac{1}{2}}^{11} \\
 & + \frac{11}{1} \frac{11619}{93536} \delta_{-5}^{10} \\
 & + \frac{105}{309} \frac{61007}{65760} \delta_{-5}^{12} \\
 & + \frac{51}{9} \frac{42611}{67680} \delta_{-4\frac{1}{2}}^9 \\
 & + \frac{84}{221} \frac{34073}{18400} \delta_{-4\frac{1}{2}}^{11} \\
 & + \frac{1}{2} \frac{03871}{41920} \delta_{-4}^{10} \\
 & - \frac{15}{387} \frac{58369}{07200} \delta_{-4}^{12} \\
 & + \frac{2317}{480} \delta_{-4}^8 \\
 & + \frac{4}{9} \frac{71539}{67680} \delta_{-3\frac{1}{2}}^9 \\
 & - \frac{24}{516} \frac{79643}{09600} \delta_{-3\frac{1}{2}}^{11} \\
 & + \frac{8197}{1920} \delta_{-3\frac{1}{2}}^7 \\
 & + \frac{357}{640} \delta_{-3}^8 \\
 & - \frac{3737}{64512} \delta_{-3}^{10} \\
 & + \frac{12}{1548} \frac{05453}{28800} \delta_{-3}^{12} \\
 & + \frac{29}{8} \delta_{-3}^6 \\
 & + \frac{1237}{1920} \delta_{-2\frac{1}{2}}^7 \\
 & - \frac{13649}{1} \frac{93536}{93536} \delta_{-2\frac{1}{2}}^9 \\
 & + \frac{72851}{73} \frac{72800}{72800} \delta_{-2\frac{1}{2}}^{11} \\
 & + \frac{3}{4} \delta_{-2}^6 \\
 & - \frac{83}{960} \delta_{-2}^8 \\
 & + \frac{1219}{96768} \delta_{-2}^{10} \\
 & - \frac{1}{774} \frac{62209}{14400} \delta_{-2}^{12} \\
 & + \frac{7}{8} \delta_{-1\frac{1}{2}}^5 \\
 & - \frac{203}{1920} \delta_{-1\frac{1}{2}}^7 \\
 & + \frac{15419}{9} \frac{67680}{67680} \delta_{-1\frac{1}{2}}^9 \\
 & - \frac{4}{1548} \frac{20529}{28800} \delta_{-1\frac{1}{2}}^{11} \\
 & + \frac{10679}{172} \frac{03200}{03200} \delta_{-1}^{12} \\
 & - \frac{1}{8} \delta_{-1}^6 \\
 & + \frac{37}{1920} \delta_{-1}^8 \\
 & - \frac{3229}{9} \frac{67680}{67680} \delta_{-1}^{10} \\
 & + \frac{10679}{172} \frac{03200}{03200} \delta_{-1\frac{1}{2}}^{11} \\
 & + \frac{1}{2} \delta_{-1}^5 \\
 & + \frac{37}{1920} \delta_{-1}^7 \\
 & - \frac{3229}{9} \frac{67680}{67680} \delta_{-1}^9 \\
 & + \frac{10679}{172} \frac{03200}{03200} \delta_{-1\frac{1}{2}}^{11} \\
 & + \frac{10879}{172} \frac{03200}{03200} \delta_{\frac{1}{2}}^{11} \\
 & + \frac{1}{8} \delta_{\frac{1}{2}}^5 \\
 & + \frac{37}{1920} \delta_{\frac{1}{2}}^7 \\
 & - \frac{3229}{9} \frac{67680}{67680} \delta_{\frac{1}{2}}^9 \\
 & + \frac{10679}{172} \frac{03200}{03200} \delta_{\frac{1}{2}}^{11} \\
 & + \frac{1}{5} \delta_{1\frac{1}{2}}^5 \\
 & + \frac{37}{1920} \delta_{1\frac{1}{2}}^7 \\
 & - \frac{3229}{9} \frac{67680}{67680} \delta_{1\frac{1}{2}}^9 \\
 & + \frac{10679}{172} \frac{03200}{03200} \delta_{1\frac{1}{2}}^{11} \\
 & + \frac{1}{3} \delta_2^6 \\
 & - \frac{37}{1920} \delta_2^8 \\
 & + \frac{3229}{9} \frac{67680}{67680} \delta_2^{10} \\
 & - \frac{10679}{172} \frac{03200}{03200} \delta_2^{12} \\
 & + \frac{7}{5} \delta_{2\frac{1}{2}}^5 \\
 & - \frac{203}{1920} \delta_{2\frac{1}{2}}^7 \\
 & + \frac{15419}{9} \frac{67680}{67680} \delta_{2\frac{1}{2}}^9 \\
 & - \frac{4}{1548} \frac{20529}{28800} \delta_{2\frac{1}{2}}^{11} \\
 & + \frac{1}{774} \frac{62209}{14400} \delta_3^{12} \\
 & - \frac{3}{4} \delta_3^6 \\
 & + \frac{83}{960} \delta_3^8 \\
 & - \frac{1219}{96768} \delta_3^{10} \\
 & + \frac{1}{774} \frac{62209}{14400} \delta_3^{12} \\
 & - \frac{23}{2} \delta_{3\frac{1}{2}}^5 \\
 & + \frac{1237}{1920} \delta_{3\frac{1}{2}}^7 \\
 & - \frac{13649}{1} \frac{93536}{93536} \delta_{3\frac{1}{2}}^9 \\
 & + \frac{72851}{73} \frac{72800}{72800} \delta_{3\frac{1}{2}}^{11} \\
 & - \frac{12}{1548} \frac{05453}{28800} \delta_4^{12} \\
 & - \frac{29}{8} \delta_4^6 \\
 & - \frac{357}{640} \delta_4^8 \\
 & + \frac{3737}{64512} \delta_4^{10} \\
 & - \frac{1548}{1548} \frac{20529}{28800} \delta_4^{11} \\
 & + \frac{15}{387} \frac{58369}{07200} \delta_5^{12} \\
 & + \frac{8197}{1920} \delta_{4\frac{1}{2}}^7 \\
 & + \frac{4}{9} \frac{71539}{67680} \delta_{4\frac{1}{2}}^9 \\
 & - \frac{1}{516} \frac{79643}{09600} \delta_{4\frac{1}{2}}^{11} \\
 & + \frac{105}{309} \frac{61007}{65760} \delta_6^{12} \\
 & - \frac{2317}{480} \delta_5^8 \\
 & + \frac{51}{9} \frac{42611}{67680} \delta_{5\frac{1}{2}}^9 \\
 & - \frac{1}{2} \frac{03871}{41920} \delta_5^{10} \\
 & + \frac{3161}{516} \frac{11237}{09600} \delta_{5\frac{1}{2}}^{11} \\
 & - \frac{11}{1} \frac{11619}{93536} \delta_6^{10} \\
 & + \frac{3161}{516} \frac{11237}{09600} \delta_{6\frac{1}{2}}^{11} \\
 & - \frac{5005}{774} \frac{69373}{14400} \delta_7^{12}
 \end{aligned}$$