

## Redetermined Crystal Structures of NiTe<sub>2</sub>, PdTe<sub>2</sub>, PtS<sub>2</sub>, PtSe<sub>2</sub>, and PtTe<sub>2</sub>

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The present study was originally started with the purpose of reinvestigating the variable  $z$  parameter in the Cd(OH)<sub>2</sub> type crystal structure of NiTe<sub>2</sub> and PdTe<sub>2</sub>. (The Cd(OH)<sub>2</sub> type structure is described in terms of space group  $P\bar{3}m1$  with metal atoms in 0,0,0 and metalloid atoms in  $\frac{2}{3}, \frac{1}{3}, z$  and  $\frac{1}{3}, \frac{2}{3}, \bar{z}$  with  $z \approx \frac{1}{4}$ .) The previous determination of the crystal structure of PdTe<sub>2</sub> by Thomassen<sup>1</sup> was based on the assumption of an ideal  $z$  value of  $\frac{1}{4}$ , whereas Tengnér<sup>2</sup> used a rotation diagram to estimate  $z = 0.25 \pm 0.01$  for NiTe<sub>2</sub>. The isostructural compounds PtS<sub>2</sub>, PtSe<sub>2</sub>, and PtTe<sub>2</sub> have recently been reinvestigated by Grønvold *et al.*<sup>3</sup> As only a qualitative criterion was used to compare observed and calculated intensities of the X-ray powder photographs, these compounds were also included in the present study.

Samples were prepared by treating the mixed powders of metal (of purity 99.9 % or better) and chalcogen (of purity 99.999+ %) in the stoichiometric ratio 1:2 in evacuated, sealed silica tubes, as previously described by Grønvold *et al.*,<sup>3</sup> Westrum *et al.*,<sup>4</sup> and Grønvold and Røst.<sup>5</sup>

All samples were crushed and X-ray powder photographs taken with filtered CuK radiation ( $\lambda(\alpha_1) = 1.54050 \text{ \AA}$ ) in cameras with 114.6 mm effective diameter and asymmetric film mounting. Photographs were also taken with strictly monochromatized CuK $\alpha_1$  radiation in a Guinier type camera with KCl as internal standard. The relative intensities of the reflections on multiple-film Debye-Scherrer photographs were determined from photometric recordings of the films. Corrections for the resolution of K $\alpha_{1,2}$  doublets were carried out according to the method of Rae and Barker.<sup>6</sup>  $F_o^2$  values were obtained by multiplication of the corrected intensities with  $(Lp \times \nu)^{-1}$ . (No corrections for absorption and temperature factors were used.) In the calculation of  $F_c$  values, atomic scattering factors were taken from *International Tables*.<sup>7</sup> The reliability index

$$R^* = \frac{\sum |F_o^2 - F_c^2|}{\sum F_o^2}$$

has been used in order to judge the accordance between the observed and calculated data. ( $F_c^2$  represents the sum of the squares of the structure factors for reflections with equal  $\sin^2\theta$ . The observed intensities of the strong, high-angle reflections ( $\theta > 75^\circ$ ) are, probably because of systematic errors, considerably lower than calculated, and these reflections are not taken into account in the calculation of  $R^*$ .)

Sets of  $F_c$  values were calculated for each compound for values of  $z$  between 0.20 and 0.28. The  $z, R^*$  curves are shown in Fig. 1. The minima of the curves indicate the most probable values of  $z$ . Unit cell dimensions,  $z$  parameters with correspond-

Table 1. Structural data for the MX<sub>2</sub> compounds  
NiTe<sub>2</sub>, PdTe<sub>2</sub>, PtS<sub>2</sub>, PtSe<sub>2</sub>, and PtTe<sub>2</sub>.

Compound	NiTe <sub>2</sub>	PdTe <sub>2</sub>	PtS <sub>2</sub>	PtSe <sub>2</sub>	PtTe <sub>2</sub>
$a$ (Å)	3.8542	4.0365	3.5432	3.7278	4.0259
$c$ (Å)	5.2604	5.1262	5.0388	5.0813	5.2209
$z$	0.254 ± 0.004	0.247 ± 0.005	0.227 ± 0.010	0.255 ± 0.003	0.254 ± 0.005
$M - 6M$ (Å)	3.854	4.037	3.543	3.728	4.026
$M - 6X$ (Å)	2.596 ± 0.011	2.652 ± 0.013	2.34 ± 0.03	2.513 ± 0.009	2.676 ± 0.013
$X - 3M$ (Å)					
$X - 3X$ (Å)	3.41 ± 0.04	3.49 ± 0.04	3.43 ± 0.08	3.29 ± 0.02	3.46 ± 0.04
$X - 3X$ (Å)	3.48 ± 0.04	3.44 ± 0.04	3.07 ± 0.08	3.37 ± 0.02	3.53 ± 0.04
$X - 6X$ (Å)	3.854	4.037	3.543	3.728	4.026

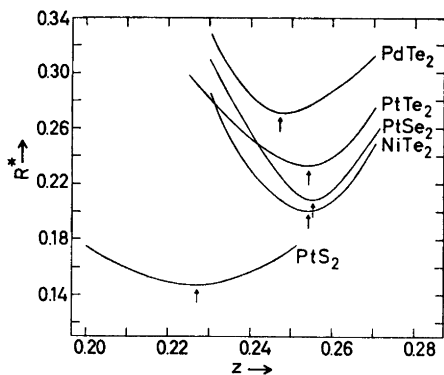


Fig. 1. The  $R^*$  versus  $z$  curves for  $\text{NiTe}_2$ ,  $\text{PdTe}_2$ ,  $\text{PtS}_2$ ,  $\text{PtSe}_2$ , and  $\text{PtTe}_2$ .

ing uncertainties estimated from the curvature of the  $z, R^*$  curves, and the shortest interatomic distances are listed in Table 1.

Some comments on the results should be appropriate:

(i) All samples used for collecting the intensity data in this study, suffered from preferred orientation of the crystals in the glass capillaries. This effect was least significant for  $\text{PtS}_2$  and most noteworthy for  $\text{PdTe}_2$ . Our first sample of  $\text{PdTe}_2$ , prepared at  $500^\circ\text{C}$ , gave  $R^* = 0.52$  as the best  $R^*$  value; the second sample, heated at  $300^\circ\text{C}$  for four days, gave the  $z, R^*$  curve shown in Fig. 1. Within the limited accuracy the position of the minimum was the same for the two sets of intensity data. This indicates that the determinations of  $z$  in Table 1 are quite reliable despite the relatively high values of  $R^*$  shown in Fig. 1.

(ii) The only compound of this investigation having a  $z$  parameter markedly different from the ideal value  $z = \frac{1}{4}$ , is  $\text{PtS}_2$ . The flatness of the  $R^*$  versus  $z$  curve for  $\text{PtS}_2$  explains why this deviation has not been observed by Grønvold *et al.*,<sup>3</sup> showing clearly that qualitative criteria may be insufficient for accurate structure determinations. (Because of the large difference in the X-ray scattering factors of Pt and S,  $R^*$  for  $\text{PtS}_2$  is rather insensitive to the variations in  $z$ ).

(iii) The reliability index  $R^{**} = \frac{\Sigma|\sqrt{F_o^2} - |\sqrt{F_c^2}|}{\Sigma|\sqrt{F_o^2}|}$  can more easily than  $R^*$  be compared with the commonly used  $R = \frac{\Sigma|F_o - |F_c||}{\Sigma|F_o|}$ . In the

minima of the  $z, R^*$  curves the corresponding values of  $R^{**}$  are 0.12, 0.13, 0.083, 0.12, and 0.11 for  $\text{NiTe}_2$ ,  $\text{PdTe}_2$ ,  $\text{PtS}_2$ ,  $\text{PtSe}_2$ , and  $\text{PtTe}_2$ , respectively. ( $R^{**}$  is 0.095 for  $\text{PtS}_2$  at  $z = 0.250$ ).

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## On the Magnetic Properties of Niobium Selenides and Tellurides

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As part of our continued studies of niobium selenides and tellurides<sup>1-4</sup> we here report the results of magnetic susceptibility measurements carried out on the existing phases  $\text{Nb}_5\text{Se}_4$ ,  $\text{Nb}_5\text{Te}_4$ ,  $\text{Nb}_{1+x}\text{Se}_2$  ( $0.00 \leq x \leq 0.29$  at  $25^\circ\text{C}$ ), " $\text{NbSe}_4$ ",  $\text{Nb}_6\text{Te}_4$ ,  $\text{Nb}_7\text{Te}_4$ ,  $\text{Nb}_2\text{Te}_2$ , and  $\text{NbTe}_4$ .

Purity of the materials and preparation of the samples have previously been described.<sup>1-4</sup> The magnetic measurements were made according to the Gouy method at temperatures between 90 and  $725^\circ\text{K}$  and at three different maximum field strengths ( $H_{\text{max}} = 4015, 4700, \text{ and } 5110 \text{ O}$ , respectively). The samples were enclosed